

$$= \left(\frac{1}{2} (\varepsilon_1 + \dots + \varepsilon_m) - \varepsilon_j \right) (h_i) a_j^* \cdot 1 \quad [287]$$

$$= \frac{1}{2} (\varepsilon_1 + \dots - \varepsilon_j + \dots + \varepsilon_m) (h_i) a_j^* \cdot 1 \quad \text{says}$$

$a_j^* \cdot 1$ has weight $\lambda - \varepsilon_j$.

In general $(a_1^*)^{k_1} \dots (a_m^*)^{k_m} \cdot 1$ has weight

$\lambda - \sum_{j=1}^m k_j \varepsilon_j$ so the set of all weights

in CM_m is $\left\{ \frac{1}{2} (\pm \varepsilon_1 \pm \varepsilon_2 \dots \pm \varepsilon_m) \right\}$ with
 indep. choices of all \pm signs, $2^m = \dim(CM)$

Def. For $i=0,1$, let $CM_m^i =$
 $\text{span} \left\{ (a_1^*)^{k_1} \dots (a_m^*)^{k_m} \cdot 1 \mid \sum_{j=1}^m k_j \equiv i \pmod{2} \right\}$.

Th. $CM_m = CM_m^0 \oplus CM_m^1$ is a direct sum decomposition of CM_m as irreducible \mathfrak{g} -modules. |288

Pf. This decomposition by the parity of the number of creation operators applied to the vacuum vector 1 , is preserved by the action of the quadratic operators $:ab:$ $\in \mathfrak{g}$. For i, j

$$:a_i^* a_i^* : \cdot 1 = a_i^* a_i^* 1 \in CM_m^0 \quad \text{and for } k \neq l,$$

$$:a_k^* a_l^* : \cdot (a_i^* a_j^* 1) = a_k^* a_l^* a_i^* a_j^* 1 =$$

$$\pm (\delta_{ki} - 1)(\delta_{kj} - 1)(\delta_{li} - 1)(\delta_{lj} - 1) a_k^* a_l^* a_i^* a_j^* 1$$

since a factor of $(a_i^*)^2 = 0$ or $(a_j^*)^2 = 0$ occurs after anti-commuting if $k \in \{i, j\}$ or $l \in \{i, j\}$.

Also, the anti-commutation relations in [289] Clifford show that the action of $:a_i a_j^*:$ for $i \neq j$ on a general basis vector of CM_n gives 0 if $k_i = 1$ or if $k_i = 0$, but when $k_i = 1$ and $k_j = 0$ it has the effect of removing a_i^* and inserting a_j^* , preserving the parity, and at most multiplying the vector by ± 1 .

Ex. $:a_1 a_2^* : \cdot (a_1^* a_3^* 1) = a_1 a_2^* a_1^* a_3^* 1 =$

$$-a_2^* a_1 a_1^* a_3^* 1 = -a_2^* (-a_1^* a_1 + 1) a_3^* 1$$

$$= a_2^* a_1^* a_1 a_3^* 1 - a_2^* a_3^* 1 =$$

$$= a_2^* a_1^* \underbrace{(-a_3^* a_1)}_{=0} 1 - a_2^* a_3^* 1 = -a_2^* a_3^* 1.$$

Similarly, the action of $:a_i a_j:$ eg $[290]$ on a general vector gives 0 if $k_i=0$ or $k_j=0$ since the annihilation operators a_i or a_j will anti-commute passed the creation operators with different subscripts and kill 1.

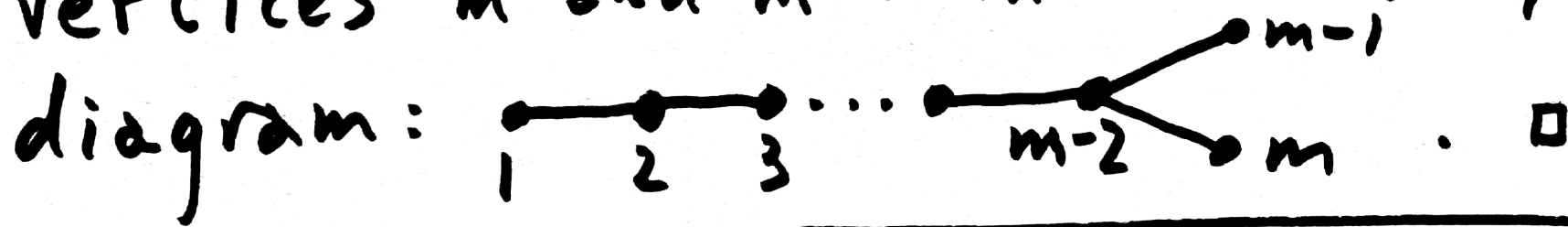
If $k_i=1=k_j$, then the action removes both a_i^* and a_j^* , leaving ± 1 times a basis vector with $k_i=0=k_j$. This preserves parity.

With the partial order on weights coming from the choice of simple roots

$\Delta = \{\alpha_1 = \epsilon_1 - \epsilon_2, \alpha_2 = \epsilon_2 - \epsilon_3, \dots, \alpha_{m-1} = \epsilon_{m-1} - \epsilon_m, \alpha_m = \epsilon_{m-1} + \epsilon_m\}$
 we find $\lambda = \frac{1}{2}(\epsilon_1 + \dots + \epsilon_m)$ is the highest weight in CM_m^0 (of the vacuum vector, 1) and

$\lambda - \epsilon_m$ is the highest weight in $(M_m)'$ (291)
 (of the vector, a_m^*).

Since $\langle \lambda, \alpha_j \rangle = \delta_{jm}$ and $\langle \lambda - \epsilon_m, \alpha_j \rangle = \delta_{j(m-1)}$
 we have $\lambda = \lambda_m$ and $\lambda - \epsilon_m = \lambda_{m-1}$ are the
 fundamental weights corresponding to the
 vertices m and $m-1$ in the D_m Dynkin



Note: The natural \mathfrak{g} -module $U \cong F^{2m}$
 with weights $\{\pm \epsilon_j \mid 1 \leq j \leq m\}$ has highest
 weight ϵ_1 (for the vector a_1) and $\lambda_1 = \epsilon_1$,
 since $\langle \epsilon_1, \alpha_j \rangle = \delta_{j1}$.

Weyl algebras:

$\text{Char}(F) \neq 2$

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Let V be a vector space over F with an anti-symm. bil. form $(\cdot, \cdot): V \times V \rightarrow F$.

Let $\text{Weyl}(V, (\cdot, \cdot))$ be the assoc. alg. with 1 generated by V with relations

$$v_1 v_2 - v_2 v_1 = (v_1, v_2) 1 \quad \forall v_1, v_2 \in V.$$

For (\cdot, \cdot) to be non-degenerate we must have $\dim(V) = 2m$ is even. Assume V has a basis $B = \{a_1, \dots, a_m, a_1^*, \dots, a_m^*\}$ s.t. the matrix of the form w.r.t. B is $M = \begin{bmatrix} 0 & I_m \\ -I_m & 0 \end{bmatrix} = J$ (Humphreys).

So the relations are:

$$a_i a_j - a_j a_i = 0 = a_i^* a_j^* - a_j^* a_i^* \quad \text{for } 1 \leq i, j \leq m$$

and

$a_i a_j^* - a_j^* a_i = \delta_{ij}$ for $1 \leq i, j \leq m$ so 1293
 $a_j^* a_i - a_i a_j^* = -\delta_{ij}$ is implied. Note that the
 case of $i=j$ in the first relations gives no
 information about a_i^2 and $(a_i^*)^2$.

Th. $\text{Weyl}_m = \text{span} \{ (a_1^*)^{k_1} \dots (a_m^*)^{k_m} a_1^{l_1} \dots a_m^{l_m} \mid 0 \leq k_i, l_i \}$
 is ∞ -dim'l.

Def. Let $\mathcal{I}_m =$ left ideal of Weyl_m gen. by
 a_1, \dots, a_m so the basis vectors with any $l_i > 0$
 are in \mathcal{I}_m .

Def. Let $\text{WM}_m = \text{Weyl}_m / \mathcal{I}_m$ which has
 a basis whose coset representatives are
 $\{ (a_1^*)^{k_1} \dots (a_m^*)^{k_m} \mathbb{1} \mid 0 \leq k_i \in \mathbb{Z} \}$ and has vacuum vector
 $\mathbb{1} = \mathbb{1} + \mathcal{I}_m$.

Th. WM_m is a Weyl $_m$ -module with left [294] mult. as the action on WM_m . The a_i act as annihilation operators, a_i^* act as creation ops.

Note. These are Bosonic operators since $a_i^n \neq 0, \forall n \geq 1, (a_i^*)^n \neq 0, \forall n \geq 1$ and generators

have commutation relations.

Let $\mathcal{U} = \text{span} \{a_i, a_i^* \mid 1 \leq i \leq m\}$ as before, but $\forall a, b \in \mathcal{U}$ let $:ab: = \frac{1}{2}(ab+ba) = :ba:$ be

the "bosonic" normally ordered product. So

$$ab - ba = (a, b) \mathbb{1} \text{ says } :ab: = \frac{1}{2}(ab + ba - (a, b) \mathbb{1})$$

$$= ab - \frac{1}{2}(a, b) \mathbb{1} = \frac{1}{2}(ba + (a, b) \mathbb{1} + ba) =$$

$$= ba + \frac{1}{2}(a, b) \mathbb{1}. \text{ Now } :a_i a_i: \neq 0 \neq :a_i^* a_i^*:$$

Def. Let $\mathfrak{g} = \text{span} \{ :ab: = \frac{1}{2}(ab+ba) \mid a, b \in U \}$. [295]

Then \mathfrak{g} is a subspace of Weyl_m with basis

$\{ :a_i a_j: \mid 1 \leq i \leq j \leq m \} \cup \{ :a_i^* a_j^* : \mid 1 \leq i \leq j \leq m \}$
 $\cup \{ :a_i a_j^* : \mid 1 \leq i, j \leq m \}$ so, counting basis vectors,

$$\text{get } \dim(\mathfrak{g}) = \frac{(m)(m+1)}{2} + \frac{m(m+1)}{2} + m^2 = 2m^2 + m$$

$$= \dim(\text{sp}(2m, F)).$$

Lemma. (Exercise) Find $[:ab:, :cd:]$ formula

as a lin. comb. from \mathfrak{g} to get \mathfrak{g} is a Lie alg.

Def. Let $h_i = :a_i a_i^* :$ for $1 \leq i \leq m$, $\mathfrak{H} = \text{span}\{h_i\}$

Cor. $\forall 1 \leq i, j \leq m$, $[h_i, h_j] = 0$. $\left. \begin{array}{l} \mathfrak{H}^* = \text{span}\{\epsilon_i\} \text{ s.t.} \\ \epsilon_i(h_j) = \delta_{ij}. \end{array} \right\}$

Cor: $[h_i, :a_j a_k:] = (\epsilon_j + \epsilon_k)(h_i) :a_j a_k:, j \leq k, \underline{296}$

$[h_i, :a_j^* a_k^*:] = -(\epsilon_j + \epsilon_k)(h_i) :a_j^* a_k^*:, j \leq k, \text{ and}$

$[h_i, :a_j a_k^*:] = (\epsilon_j - \epsilon_k)(h_i) :a_j a_k^*:$ for $1 \leq j, k \leq m,$

so $\mathfrak{g} = H \oplus \bigoplus_{\alpha \in \Phi} \mathfrak{g}_\alpha$ where $\Phi = \{\pm \epsilon_j \pm \epsilon_k \mid 1 \leq j \leq k \leq m\}$

but $0 \notin \Phi,$ and

$\mathfrak{g}_\alpha = F :a_j a_k:$ for $\alpha = \epsilon_j + \epsilon_k$ with $1 \leq j \leq k \leq m,$

$\mathfrak{g}_\alpha = F :a_j^* a_k^*:$ for $\alpha = -\epsilon_j - \epsilon_k$ with $1 \leq j \leq k \leq m,$

$\mathfrak{g}_\alpha = F :a_j a_k^*:$ for $\alpha = \epsilon_j - \epsilon_k$ with $1 \leq j \neq k \leq m.$

Th: $\mathfrak{g} \cong \mathfrak{sp}(2m, F)$ where the isomorphism can

be made to match corresponding root vectors in \mathfrak{g} with specific matrices in the $\mathfrak{sp}(2m, F)$ basis.

Note: The representation of \mathfrak{g} on WM_m (297) is ∞ -dim'l, but $1 \in WM_m$ is a highest weight vector w.r.t. simple roots (type C_m)

$$\Delta = \{\alpha_1 = \epsilon_1 - \epsilon_2, \alpha_2 = \epsilon_2 - \epsilon_3, \dots, \alpha_{m-1} = \epsilon_{m-1} - \epsilon_m, \alpha_m = 2\epsilon_m\}$$

and we can compute the weight λ as follows.

$$h_i \cdot 1 = - : a_i a_i^* : 1 = -(a_i^* a_i + \frac{1}{2}) 1 = -\frac{1}{2} \cdot 1 \quad \text{for } 1 \leq i \leq m$$

would say $\lambda = -\frac{1}{2}(\epsilon_1 + \dots + \epsilon_m)$ so $\lambda(h_i) = -\frac{1}{2}$ for $1 \leq i \leq m$.

This might need modification if our def. of h_i is off by a scalar! Done.

Th: $WM_m = WM_m^0 \oplus WM_m^1$ is a decomp. into irred. \mathfrak{g} -modules, where WM_m^i for $i=0,1$ is a parity decomposition as before.

We may also obtain the natural representation of g on $U \cong F^{2m}$ as follows. 298

Define left action of g on U by

$$(:ab:) \cdot c = [(:ab:), c] = [ab, c] = abc - cab$$

$$= abc - acb$$

$$+ acb - cab$$

$$= a(bc - cb) + (ac - ca)b$$

$$= (b, c)a + (a, c)b \in U.$$

Check this gives a Lie alg. rep'n of g on U as on page 283.

Exercise: Find $h_i \cdot a_j$ and $h_i \cdot a_j^*$ to get weights of these vectors. If we want $a_j \leftrightarrow \epsilon_j$ and $a_j^* \leftrightarrow -\epsilon_j$ we need $h_i = -:a_i a_i^*:$