# Functional determinants for the Dirac equation with mixed pseudodifferential boundary conditions over finite cylinders 

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(Received 28 July 2004; accepted 27 November 2004; published online 28 March 2005)


#### Abstract

In this note, we explicitly compute the functional determinant of a Dirac Laplacian with nonlocal pseudodifferential boundary conditions over a finite cylinder in terms of the $\zeta$-function of the Dirac operator on the cross section and the pseudodifferential operators defining the boundary conditions. In particular, this result reduces to our previous formula [J. Phys. A 37, 7381 (2004)] for the special case of generalized Atiyah-Patodi-Singer conditions. To prove our main result, we use the gluing and comparison formulas established by the present authors in Refs 14 and 15.


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## I. INTRODUCTION

Recent advances in quantum field theory have necessitated the explicit evaluation of functional determinants of Dirac operators over a variety of space-time configurations. In fact, at the one-loop order, any such theory can be reduced to the theory of determinants. We refer the reader to the works of Dowker and Critchley ${ }^{1}$ and Hawking. ${ }^{2}$ See also Elizalde et al., ${ }^{3}$ Kirsten, ${ }^{4}$ and Scott and Wojciechowski, ${ }^{5}$ for recent reviews. Because of their increasingly important role in mathematical physics, over the past several years there has been intense research to compute functional $\zeta$-determinants of Dirac Laplacians. Of great significance is the Dirac Laplacian with spectral pseudodifferential boundary conditions; the Atiyah-Patodi-Singer (henceforth APS) boundary conditions being the most well-known example. Such boundary conditions arise in, for instance, one-loop quantum cosmology, ${ }^{6-8}$ spectral branes, ${ }^{9}$ and the study of Dirac fields in the background of a magnetic flux. ${ }^{10}$

However, only recently was the open problem of explicitly computing the $\zeta$-determinant of a Dirac Laplacian with APS conditions over a finite cylinder solved. ${ }^{11}$ One reason this problem withstood the efforts of existing mathematical methods is that it is impossible to find explicit formulas for the eigenvalues of such a Dirac operator. For this reason, we had to attack the problem using the method of adiabatic decomposition pioneered by Douglas and Wojciechowski ${ }^{12}$ for the eta invariant and by the second author and Wojciechowski ${ }^{13}$ for the $\zeta$-determinant. The purpose of this current paper is twofold. First, we extend the result of Ref. 11 to a general class of pseudodifferential conditions that generalize the APS condition up to operators of arbitrary finite rank. To compute the $\zeta$-determinant in this generalized framework, which in some sense possesses eigenvalues that are even more enigmatic, we use the gluing and comparison formulas for $\zeta$-determinants proved by the authors in Refs. 14 and 15 to break up this general framework into tangible parts which can be explicitly computed. The second purpose of this paper is to elucidate

[^0]the effectiveness of these gluing and comparison formulas to compute $\zeta$-determinants that have eluded explicit evaluations due to the perplexity of their eigenvalues. This also exemplifies the aim of gluing and comparison formulas: Breaking up complex problems into simpler more tractable ones.

We now describe our setup. Fix $R>0$ and let $\mathcal{D}: C^{\infty}\left(N_{R}, S\right) \rightarrow C^{\infty}\left(N_{R}, S\right)$ be a Dirac type operator where $N_{R}=[-R, R] \times Y$ is a finite cylinder with $R>0, Y$ a closed compact Riemannian manifold (of arbitrary dimension), and $S$ a Clifford bundle over $N_{R}$. We assume that $\mathcal{D}$ is of product form

$$
\begin{equation*}
\mathcal{D}=G\left(\partial_{u}+D_{Y}\right), \tag{1.1}
\end{equation*}
$$

where $G$ is a unitary bundle isomorphism of $S$ and $D_{Y}$ is a Dirac operator acting on $C^{\infty}(Y, S)$ such that $G^{2}=-\mathrm{Id}$ and $G D_{Y}=-D_{Y} G$. Furthermore, we assume that

$$
\begin{equation*}
\operatorname{dim}\left(\operatorname{ker}(G+i) \cap \operatorname{ker}\left(D_{Y}\right)\right)=\operatorname{dim}\left(\operatorname{ker}(G-i) \cap \operatorname{ker}\left(D_{Y}\right)\right) \tag{1.2}
\end{equation*}
$$

Let $\Pi_{+}, \Pi_{-}$, and $\Pi_{0}$ denote the orthogonal projections onto the positive, negative, and zero eigenspaces of $D_{Y}$. Since $N_{R}$ has boundaries, we have to impose boundary conditions. Let $\operatorname{Gr}_{\infty}^{*}\left(D_{Y}\right)$ denote the space of pairs $\left(\mathcal{P}_{1}, \mathcal{P}_{2}\right)$, where $\mathcal{P}_{1}$ and $\mathcal{P}_{2}$ are orthogonal pseudodifferential projections on $L^{2}(Y, S)$ such that

$$
\mathcal{P}_{1}-\Pi_{+}, \quad \mathcal{P}_{2}-\Pi_{-} \text {are smoothing operators, }
$$

and for $i=1,2$,

$$
\mathcal{P}_{i} G=G\left(\operatorname{Id}-\mathcal{P}_{i}\right), \quad D_{Y} \mathcal{P}_{i}=\mathcal{P}_{i} D_{Y} .
$$

An important class of such boundary conditions are the renowned generalized APS spectral conditions, ${ }^{16}$ which are defined as follows. Let $\sigma_{1}, \sigma_{2}$ be involutions (that is, $\sigma_{i}^{2}=\mathrm{Id}$ ) over $\operatorname{ker}\left(D_{Y}\right)$ such that $\sigma_{1} G=-G \sigma_{1}$ and $\sigma_{2} G=-G \sigma_{2}$. Note that such involutions exist because of the assumption (1.2). Then

$$
\begin{equation*}
\Pi_{+}^{\sigma_{1}}:=\Pi_{+}+\frac{1+\sigma_{1}}{2} \Pi_{0}, \quad \Pi_{-}^{\sigma_{2}}:=\Pi_{-}+\frac{1+\sigma_{2}}{2} \Pi_{0} \tag{1.3}
\end{equation*}
$$

are called generalized APS spectral projections, and $\left(\Pi_{+}^{\sigma_{1}}, \Pi_{-}^{\sigma_{2}}\right) \in \operatorname{Gr}_{\infty}^{*}\left(D_{Y}\right)$. These generalized APS boundary conditions were considered in our paper (Ref. 11), but elements of $\operatorname{Gr}_{\infty}^{*}\left(D_{Y}\right)$ are much more general and can differ from APS projections by operators of arbitrary finite rank. Let $P$ $=\left(\mathcal{P}_{1}, \mathcal{P}_{2}\right) \in \operatorname{Gr}_{\infty}^{*}\left(D_{Y}\right)$ and impose boundary conditions for $\mathcal{D}$ at $\{-R\} \times Y$ and $\{R\} \times Y$ via

$$
\mathcal{P}_{1} \quad \text { at }\{-R\} \times Y, \quad \mathcal{P}_{2} \quad \text { at }\{R\} \times Y .
$$

We denote by $\mathcal{D}_{P}$ the resulting operator with these boundary conditions, that is,

$$
\begin{equation*}
\mathcal{D}_{P}:=\mathcal{D}: \operatorname{dom}\left(\mathcal{D}_{P}\right) \rightarrow L^{2}\left(N_{R}, S\right), \tag{1.4}
\end{equation*}
$$

where

$$
\operatorname{dom}\left(\mathcal{D}_{P}\right):=\left\{\phi \in H^{1}\left(N_{R}, S\right) \mid \mathcal{P}_{1}\left(\left.\phi\right|_{u=-R}\right)=0, \quad \mathcal{P}_{2}\left(\left.\phi\right|_{u=R}\right)=0\right\} .
$$

By the fundamental work of Seeley, ${ }^{17,18}$ the spectrum of the Dirac operator $\mathcal{D}_{P}$ consists of discrete real eigenvalues $\left\{\lambda_{k}\right\}$. The $\zeta$-function of $\mathcal{D}_{P}^{2}$ is defined by

$$
\zeta_{\mathcal{D}_{P}^{2}}(s)=\sum_{\lambda_{k} \neq 0} \lambda_{k}^{-2 s},
$$

which is a priori defined for $\mathfrak{R}(s) \gg 0$, and by the work of Grubb ${ }^{19,20}$ and Wojciechowski ${ }^{21}$ (cf. Lei ${ }^{22}$ and Loya and Park ${ }^{15}$ ), has a meromorphic extension to C with 0 as a regular point. The $\zeta$-determinant of $\mathcal{D}_{P}^{2}$ is defined by

$$
\operatorname{det}_{\zeta} \mathcal{D}_{P}^{2}:=\exp \left(-\zeta_{\mathcal{D}_{P}^{2}}^{\prime}(0)\right)
$$

This definition first appeared in Ray and Singer's seminal paper ${ }^{23}$ on the analytic torsion. Since we imposed nonlocal pseudodifferential boundary conditions, it is impossible to compute the eigenvalues $\left\{\lambda_{k}\right\}$ of $\mathcal{D}_{P}$ explicitly, so there is no direct way to compute the $\zeta$-determinant $\operatorname{det}_{\xi} \mathcal{D}_{P}^{2}$ from the eigenvalues. However, we shall give two derivations of the formula for $\operatorname{det}_{\zeta} \mathcal{D}_{P}^{2}$ :
(1) Using the gluing formula proved in Ref. 14.
(2) Using the comparison/relative invariant formula proved in Ref. 15.

See Sec. II for more on these results. The formula for $\operatorname{det}_{\zeta} \mathcal{D}_{P}^{2}$ is described as follows. Since the $\mathcal{P}_{i}$ 's are orthogonal projectors such that $\mathcal{P}_{i} G=G\left(\operatorname{Id}-\mathcal{P}_{i}\right)$ by assumption, with respect to the decomposition

$$
\begin{equation*}
L^{2}(Y, S)=L^{2}\left(Y, S^{+}\right) \oplus L^{2}\left(Y, S^{-}\right) \tag{1.5}
\end{equation*}
$$

with $S^{ \pm}$denoting the $( \pm i)$-eigenspaces of $G$ in $S$ (recall that $G^{2}=-\mathrm{Id}$ ), we can write

$$
\mathcal{P}_{i}=\frac{1}{2}\left(\begin{array}{cc}
\operatorname{Id} & \kappa_{i}^{-1} \\
\kappa_{i} & \mathrm{Id}
\end{array}\right),
$$

where $\kappa_{i}: L^{2}\left(Y, S^{+}\right) \rightarrow L^{2}\left(Y, S^{-}\right)$are corresponding isometries. In particular,

$$
U_{P}:=-\kappa_{1} \kappa_{2}^{-1}: L^{2}\left(Y, S^{-}\right) \rightarrow L^{2}\left(Y, S^{-}\right)
$$

is a unitary operator. Let $\hat{U}_{P}$ denote the restriction of $U_{P}$ to the orthogonal complement of its $(-1)$-eigenspace. Let $W:=\operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{2}\right)$. Then our assumptions on $\left(\mathcal{P}_{1}, \mathcal{P}_{2}\right)$ imply that $W$ is a finite-dimensional vector space and that $D_{Y}: \operatorname{Im}\left(\mathcal{P}_{i}\right) \rightarrow \operatorname{Im}\left(\mathcal{P}_{i}\right)$ and $D_{Y}: \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{i}\right) \rightarrow \operatorname{Im}(\operatorname{Id}$ $-\mathcal{P}_{i}$ ). We define a map $T_{P}$ over $W$ by

$$
T_{P}:=\left\{\begin{array}{l}
\frac{\sinh \left(2 R D_{Y}\right)}{D_{Y}} \quad \text { on } W \cap \operatorname{ker}\left(D_{Y}\right)^{\perp}, \\
2 R \quad \text { on } W \cap \operatorname{ker}\left(D_{Y}\right) .
\end{array}\right.
$$

We also define $T_{\mathcal{P}_{1}}$ and $T_{\mathcal{P}_{2}}$ over the finite-dimensional vector spaces $\operatorname{Im}\left(\Pi_{+}\right) \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right)$ and $\operatorname{Im}\left(\Pi_{+}\right) \cap \operatorname{Im}\left(\mathcal{P}_{2}\right)$, respectively, by

$$
T_{\mathcal{P}_{1}}:=e^{4 R D_{Y}}, \quad T_{\mathcal{P}_{2}}:=e^{4 R D_{Y}} .
$$

The following theorem is the main result of this note.
Theorem 1.1. The following equality holds:

$$
\operatorname{det}_{\zeta} \mathcal{D}_{P}^{2}=e^{2 C R_{2}{ }^{\zeta} D_{Y}^{2}(0)+h_{Y}} \frac{\left(\operatorname{det} T_{P}\right)^{2}}{\left(\operatorname{det} T_{\mathcal{P}_{1}}\right)^{2}\left(\operatorname{det} T_{\mathcal{P}_{2}}\right)^{2}} \cdot \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{U}_{P}+\hat{U}_{P}^{-1}}{4}\right),
$$

where $C=\left(\Gamma(s)^{-1} \zeta_{D_{Y}^{2}}(s-1 / 2)\right)^{\prime}(0)$ with $\zeta_{D_{Y}^{2}}(s)$ the $\zeta$-function of $D_{Y}^{2}, h_{Y}=\operatorname{dim} \operatorname{ker}\left(D_{Y}\right)$ and $\operatorname{det}_{F}$ denotes the Fredholm determinant.

More explicitly, if we let $E_{\mu}$ denote the eigenspace of $D_{Y}$ associated to the eigenvalue $\mu$ $\in \operatorname{spec}\left(D_{Y}\right)$, then we can write

$$
\begin{align*}
\operatorname{det}_{\zeta} \mathcal{D}_{P}^{2}= & e^{2 C R_{2}{ }_{2}^{D_{Y}^{2}(0)+h_{Y}}}\left(\begin{array}{l}
\substack{\mu>0 \\
E_{\mu} \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \neq 0}
\end{array} e^{-4 \mu R}\right) \cdot\left(\begin{array}{|}
\substack{\mu>0 \\
E_{\mu} \cap \operatorname{Im}\left(\mathcal{P}_{2}\right) \neq 0} \\
& \left.e^{-4 \mu R}\right) \\
& \cdot\left(\prod_{\substack{\mu \in \operatorname{spec}\left(D_{Y}\right) \\
E_{\mu} \cap W \neq 0}} \frac{\sinh ^{2}(2 \mu R)}{\mu^{2}}\right) \cdot \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{U}_{P}+\hat{U}_{P}^{-1}}{4}\right)
\end{array},\right.
\end{align*}
$$

where if $\mu=0$ in the product in the second line, we replace $\left[\sinh ^{2}(2 \mu R)\right] / \mu^{2}$ by its limit as $\mu$ $\rightarrow 0$, that is, $(2 R)^{2}$. In particular, when $\mathcal{P}_{1}=\Pi_{+}^{\sigma_{1}}$ and $\mathcal{P}_{2}=\Pi_{-}^{\sigma_{2}}$, the generalized APS spectral projectors in (1.3), then Theorem 1.1 reduces to the main result of Ref. 11.

$$
\begin{equation*}
\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma_{1}, \Pi_{-}^{\sigma_{2}}}}^{2}=e^{2 C R_{2} \zeta_{D_{Y}^{2}}^{2}(0)+h_{Y}}(2 R)^{2 h} \operatorname{det}^{*}\left(\frac{2 \operatorname{Id}-\left(\sigma_{1} \sigma_{2}\right)_{-}-\left(\sigma_{1} \sigma_{2}\right)_{-}^{-1}}{4}\right) \tag{1.7}
\end{equation*}
$$

where $\left(\sigma_{1} \sigma_{2}\right)_{-}$is the restriction of the unitary map $\sigma_{1} \sigma_{2}$ to $L^{2}\left(Y, S^{-}\right) \cap \operatorname{ker}\left(D_{Y}\right), h$ is the number of $(+1)$-eigenvalues of $\left(\sigma_{1} \sigma_{2}\right)_{-}$, and where $\operatorname{det}^{*}(L)$ denotes the determinant of $\left(\left.L\right|_{\left.\operatorname{ker}(L)^{\perp}\right)}\right.$ for an operator $L$ over a finite-dimensional vector space.

The structure of this paper is as follows. In Sec. II, we review the gluing formula from Ref. 14 and the comparison/relative invariant formula from Ref. 15, which we shall use in the subsequent sections. In Sec. III, we derive new formulas for $\zeta$-determinant ratios of Dirac Laplacians with boundary conditions of special types. Finally, in Sec. IV we combine these special $\zeta$-determinant ratios and the gluing and comparison formulas from Refs. 14 and 15 to derive our main Theorem 1.1.

## II. THE GLUING AND COMPARISON FORMULAS FROM REFS. 14 and 15

In this section we review the gluing formula from Ref. 14 and the comparison/relative invariant formula from Ref. 15.

Let $\mathcal{D}$ be a Dirac type operator acting on $C^{\infty}(M, S)$ where $M$ is a closed compact Riemannian manifold of arbitrary dimension and $S$ is a Clifford bundle over $M$. Suppose that $M=M_{-} \cup M_{+}$is partitioned into a union of manifolds with a common boundary $Y=\partial M_{-}=\partial M_{+}$. We assume that all geometric structures are of product type over a tubular neighborhood $N$ of $Y$ where $\mathcal{D}$ takes the product form (1.1). By restriction of $\mathcal{D}$, we obtain Dirac type operators $\mathcal{D}_{ \pm}$over $M_{ \pm}$. We impose the boundary conditions given by the orthogonalized Calderón projectors $\mathcal{C}_{ \pm}$for $\mathcal{D}_{ \pm}$and we denote by $\mathcal{D}_{\mathcal{C}_{ \pm}}$the resulting operators,

$$
\mathcal{D}_{\mathcal{C}_{ \pm}}=\mathcal{D}_{ \pm}: \operatorname{dom}\left(\mathcal{D}_{\mathcal{C}_{ \pm}}\right):=\left\{\phi \in H^{1}\left(M_{ \pm}, S\right) \mid \mathcal{C}_{ \pm}\left(\left.\phi\right|_{Y}\right)=0\right\} \rightarrow L^{2}\left(M_{ \pm}, S\right) .
$$

Here, we recall that the Calderón projectors $\mathcal{C}_{ \pm}$are the projectors defined intrinsically as the unique orthogonal projectors onto the closures in $L^{2}(Y, S)$ of the infinite-dimensional Cauchy data spaces of $\mathcal{D}_{ \pm}$:

$$
\left\{\left.\phi\right|_{Y} \mid \phi \in C^{\infty}\left(M_{ \pm}, S\right), \quad \mathcal{D}_{ \pm} \phi=0\right\} \subset C^{\infty}(Y, S)
$$

The gluing problem for the $\zeta$-determinant is to describe the "defect"

$$
\frac{\operatorname{det}_{\zeta} \mathcal{D}^{2}}{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{C}_{+}}^{2} \cdot \operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{C}_{-}}^{2}}=?
$$

in terms of recognizable data. To describe the solution in Ref. 14, we need to introduce some notations. With respect to the decomposition as in (1.5), the Calderón projectors $\mathcal{C}_{ \pm}$have the matrix forms

$$
\mathcal{C}_{ \pm}=\frac{1}{2}\left(\begin{array}{cc}
\mathrm{Id} & \kappa_{ \pm}^{-1}  \tag{2.1}\\
\kappa_{ \pm} & \mathrm{Id}
\end{array}\right),
$$

where the maps $\kappa_{ \pm}: L^{2}\left(Y, S^{+}\right) \rightarrow L^{2}\left(Y, S^{-}\right)$are corresponding isometries, so that $\mathcal{U}:=-\kappa_{-} \kappa_{+}^{-1}$ is a unitary operator over $L^{2}\left(Y, S^{-}\right)$. Furthermore, $\mathcal{U}$ is of Fredholm determinant class. We denote by $\hat{\mathcal{U}}$ the restriction of $\mathcal{U}$ to the orthogonal complement of its $(-1)$-eigenspace. We also put

$$
\mathcal{L}:=\sum_{k=1}^{h_{M}} \gamma U_{k} \otimes \gamma U_{k}: \gamma \operatorname{ker}(\mathcal{D}) \rightarrow \gamma \operatorname{ker}(\mathcal{D})
$$

where $h_{M}=\operatorname{dim} \operatorname{ker}(\mathcal{D}), \gamma$ is the restriction map from $M$ to $Y$, and $\left\{U_{k}\right\}$ is an orthonormal basis of $\operatorname{ker}(\mathcal{D})$. Then $\mathcal{L}$ is a positive operator on the finite-dimensional vector space

$$
\gamma \operatorname{ker}(\mathcal{D}) \equiv \operatorname{Im}\left(\mathcal{C}_{-}\right) \cap \operatorname{Im}\left(\mathcal{C}_{+}\right)
$$

We now have all the ingredients to state the following gluing formula: ${ }^{14}$

$$
\begin{equation*}
\frac{\operatorname{det}_{\zeta} \mathcal{D}^{2}}{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{C}_{-}}^{2} \cdot \operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{C}_{+}}^{2}}=2^{-\zeta_{D_{Y}^{2}}^{2}(0)-h_{Y}}(\operatorname{det} \mathcal{L})^{-2} \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{\mathcal{U}}+\hat{\mathcal{U}}^{-1}}{4}\right) \tag{2.2}
\end{equation*}
$$

where $h_{Y}=\operatorname{dim} \operatorname{ker}\left(D_{Y}\right)$ and $\operatorname{det}_{F}$ denotes the Fredholm determinant. There is a similar formula for manifolds with cylindrical ends. ${ }^{24}$

We now explain the comparison/relative invariant formula proved in Ref. 15 for ( $M_{-}, \mathcal{D}_{-}$). To this end, we consider the space $\operatorname{Gr}_{\infty}^{*}\left(\mathcal{D}_{-}\right)$, which consists of orthogonal projections $\mathcal{P}$ such that $\mathcal{P}-\mathcal{C}_{-}$are smoothing operators and $G \mathcal{P}=(\operatorname{Id}-\mathcal{P}) G$. Let us fix $\mathcal{P} \in \operatorname{Gr}_{\infty}^{*}\left(\mathcal{D}_{-}\right)$and let $\kappa_{\mathcal{P}}: L^{2}\left(Y, S^{+}\right)$ $\rightarrow L^{2}\left(Y, S^{-}\right)$be the map that determines $\mathcal{P}$ as $\kappa_{ \pm}$does $\mathcal{C}_{ \pm}$in (2.1). Let $\mathcal{D}_{\mathcal{P}}$ denote the operator $\mathcal{D}_{-}$ on $M_{-}$with the boundary condition given by $\mathcal{P}$. Let $P_{W}$ be the orthogonal projection of $L^{2}(Y, S)$ onto the finite-dimensional vector space

$$
W:=\gamma \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}}\right) \equiv \operatorname{Im}\left(\mathcal{C}_{-}\right) \cap \operatorname{Im}(\operatorname{Id}-\mathcal{P}) .
$$

Then we introduce a linear map

$$
\begin{equation*}
L:=-P_{W} G \mathcal{R}^{-1} G P_{W}: \gamma \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}}\right) \rightarrow \gamma \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}}\right), \tag{2.3}
\end{equation*}
$$

where $\mathcal{R}$ is the sum of the Dirichlet to Neumann maps on an extension of $M_{-}$defined as follows. Fix any invertible extension $\tilde{\mathcal{D}}$ of $\mathcal{D}$ to a manifold $\tilde{M}$ that contains $M_{-}$. (The double of $\mathcal{D}$ would do nicely.) Then for any $\varphi \in C^{\infty}(Y, S)$, there are unique $\phi_{1} \in C^{\infty}\left(M_{-}, S\right)$ and $\phi_{2} \in C^{\infty}\left(\tilde{M} \backslash M_{-}, S\right)$ that are continuous at $Y$ with value $\varphi$ such that $\widetilde{\mathcal{D}}^{2} \phi_{i}=0, i=1,2$, off of $Y$. Then

$$
\begin{equation*}
\mathcal{R} \varphi:=\left.\partial_{u} \phi_{1}\right|_{Y}-\left.\partial_{u} \phi_{2}\right|_{Y} . \tag{2.4}
\end{equation*}
$$

In Ref. 15, we prove that $L$ is a positive operator so that $\operatorname{det} L$ is a positive real number. Now the main result of Ref. 15 states that

$$
\begin{equation*}
\frac{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}}^{2}}{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{C}_{-}}^{2}}=(\operatorname{det} L)^{2} \cdot \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{U}+\hat{U}^{-1}}{4}\right) \tag{2.5}
\end{equation*}
$$

where $\hat{U}$ is the restriction of $U:=\kappa_{-} \kappa_{\mathcal{P}}^{-1}$ to the orthogonal complement of its ( -1 )-eigenspace. The formula (2.5) extends the work of Scott ${ }^{25}$ for the invertible case (cf. Scott and Wojciechowski ${ }^{26}$ ), and has recently been further extended to noncompact manifolds whose boundaries are manifolds with multi-cylindrical ends. ${ }^{27}$


FIG. 1. Cutting $N_{R}$ at $r$ into two pieces.

## III. THE $\zeta$-DETERMINANT FOR SPECIAL BOUNDARY CONDITIONS

If $\left(\mathcal{P}_{1}, \mathcal{P}_{2}\right) \in \operatorname{Gr}_{\infty}^{*}\left(D_{Y}\right)$, then for the sake of clarity we shall denote the operator $\mathcal{D}_{P}$ in the introduction with these boundary conditions by $\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}$. Thus, over $N_{R}=[-R, R] \times Y$ we impose boundary conditions at $\{-R\} \times Y$ and $\{R\} \times Y$ via

$$
\mathcal{P}_{1} \text { at }\{-R\} \times Y, \quad \mathcal{P}_{2} \text { at }\{R\} \times Y,
$$

and $\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}$ is the operator with domain

$$
\begin{equation*}
\operatorname{dom}\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}\right):=\left\{\phi \in H^{1}\left(N_{R}, S\right) \mid \mathcal{P}_{1}\left(\left.\phi\right|_{u=-R}\right)=0, \quad \mathcal{P}_{2}\left(\left.\phi\right|_{u=R}\right)=0\right\} . \tag{3.1}
\end{equation*}
$$

Let $\mathcal{P}$ be a projection on $L^{2}(Y, S)$ with $(\mathcal{P}, \operatorname{Id}-\mathcal{P}) \in \operatorname{Gr}_{\infty}^{*}\left(D_{Y}\right)$. By definition of $\operatorname{Gr}_{\infty}^{*}\left(D_{Y}\right)$, the image of $\mathcal{P} \Pi_{0}$ is a Lagrangian subspace in $\operatorname{Im}\left(\Pi_{0}\right)=\operatorname{ker}\left(D_{Y}\right)$. Let $\sigma$ be the involution over $\operatorname{Im}\left(\Pi_{0}\right)$ such that $[(1+\sigma) / 2] \Pi_{0}=\mathcal{P} \Pi_{0}$. Recalling $\Pi_{+}^{\sigma}:=\Pi_{+}+[(\mathrm{Id}+\sigma) / 2] \Pi_{0}$, it follows that $\left(\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}\right)$ $\in \operatorname{Gr}_{\infty}^{*}\left(D_{Y}\right)$. Recall that we can write

$$
\mathcal{P}=\frac{1}{2}\left(\begin{array}{cc}
\mathrm{Id} & \kappa_{\mathcal{P}}^{-1} \\
\kappa_{\mathcal{P}} & \mathrm{Id}
\end{array}\right), \quad \Pi_{+}^{\sigma}=\frac{1}{2}\left(\begin{array}{cc}
\mathrm{Id} & \kappa_{\sigma}^{-1} \\
\kappa_{\sigma} & \mathrm{Id}
\end{array}\right)
$$

for corresponding isometries $\kappa_{\mathcal{P}}, \kappa_{\sigma}: L^{2}\left(Y, S^{+}\right) \rightarrow L^{2}\left(Y, S^{-}\right)$, and define $\hat{U}_{\mathcal{P}}$ as the restriction of $U_{\mathcal{P}}:=\kappa_{\mathcal{P}} \kappa_{\sigma}^{-1}$ over $L^{2}\left(Y, S^{-}\right)$to the orthogonal complement of its ( -1 )-eigenspace. Note that $-\kappa_{\mathcal{P}}$ and $-\kappa_{\sigma}$ are the isometries corresponding to $\mathrm{Id}-\mathcal{P}$ and $\mathrm{Id}-\Pi_{+}^{\sigma}$, respectively. We begin by computing the following ratio.

Lemma 3.1: We have

$$
\frac{\operatorname{det}_{\xi} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{2}}{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{P}}}=\left(\prod_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}(\mathrm{I}-\mathcal{P}) \neq 0}} \frac{e^{4 \mu \sinh ^{2}(2 \mu R)}}{\mu^{2}}\right) \cdot \operatorname{det}_{F}\left(\frac{2 \mathrm{Id}+\hat{U}_{\mathcal{P}}+\hat{U}_{\mathcal{P}}^{-1}}{4}\right),
$$

where $E_{\mu}$ denotes the eigenspace of $D_{Y}$ associated to the eigenvalue $\mu \in \operatorname{spec}\left(D_{Y}\right)$.
Proof: We give two proofs, first using the gluing formula (2.2) then using the comparison formula (2.5), in order to demonstrate the effectiveness of these formulas.

Gluing proof of Lemma 3.1: Let us decompose $N_{R}$ into two parts $[-R, r] \times Y$ and $[r, R] \times Y$ as shown in Fig. 1. Then the restrictions of $\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}$ over the decomposed parts $[-R, r] \times Y$ and $[r, R] \times Y$ define two Dirac type operators with boundary conditions given by $\mathcal{P}$ at $\{-R\} \times Y$ and Id $-\Pi_{+}^{\sigma}$ at $\{R\} \times Y$. It is easy to check that the Calderón projections of these two operators are, respectively, just $\mathcal{C}_{-}=\mathrm{Id}-\mathcal{P}$ and $\mathcal{C}_{+}=\Pi_{+}^{\sigma}$ over $\{r\} \times Y$. We denote the operators over the decomposed parts $[-R, r] \times Y$ and $[r, R] \times Y$, with the new boundary conditions given by these Calderón projectors over $\{r\} \times Y$, by $\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{P}}(r)$ and $\mathcal{D}_{\Pi_{+}^{\sigma} \mathrm{Id}-\Pi_{+}^{( }}(r)$, respectively. Applying the gluing formula (2.2) to this situation, we obtain

$$
\begin{equation*}
\frac{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{2}}{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{P}}^{2}(r) \cdot \operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma}}^{2} \mathrm{Id}-\Pi_{+}^{\sigma}(r)}=2^{-\zeta_{D_{Y}^{2}}^{2}(0)-h_{Y}}\left(\operatorname{det} \mathcal{L}_{r}\right)^{-2} \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{U}_{\mathcal{P}}+\hat{U}_{\mathcal{P}}^{-1}}{4}\right) \tag{3.2}
\end{equation*}
$$

where $\zeta_{D_{Y}^{2}}(s)$ is the $\zeta$-function of $D_{Y}^{2}$ and $h_{Y}=\operatorname{dim} \operatorname{ker}\left(D_{Y}\right)$, where we used that the $\hat{\mathcal{U}}$ in (2.2) for this situation is $-\left(-\kappa_{\mathcal{P}}\right)\left(\kappa_{\sigma}^{-1}\right)=\kappa_{\mathcal{P}} \kappa_{\sigma}^{-1}=: \hat{U}_{\mathcal{P}}$, and where

$$
\mathcal{L}_{r}:=\sum_{k=1}^{h_{\mathcal{P}}} \gamma_{r} U_{k} \otimes \gamma_{r} U_{k}: \gamma_{r} \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}\right) \rightarrow \gamma_{r} \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}\right)
$$

where $h_{\mathcal{P}}=\operatorname{dim} \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}\right), \gamma_{r}$ is the restriction map from $N_{R}$ to $\{r\} \times Y$, and $\left\{U_{k}\right\}$ is an orthonormal basis of $\operatorname{ker}\left(\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{+}\right)$. From the main result of Ref. 11 [that is, the formula (1.7)], we have

$$
\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma} \text { Id }-\Pi_{+}^{\sigma}}^{2}(r)=e^{C(R-r)} 2^{\zeta_{D} D_{Y}^{2}(0)+h_{Y}},
$$

where $C=\left(\Gamma(s)^{-1} \zeta_{D_{Y}^{2}}(s-1 / 2)\right)^{\prime}(0)$. Thus, $\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma}}^{2}{ }_{\text {Id }-\Pi_{+}^{\sigma}}(r) \rightarrow 2^{\zeta_{D_{Y}^{2}}^{2}(0)+h_{Y}}$ as $r \rightarrow R$, so taking $r \rightarrow R$ in (3.2), we see that

$$
\frac{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{2}}{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{P}}^{2}}=\left(\operatorname{det} \mathcal{L}_{R}\right)^{-2} \operatorname{det}_{F}\left(\frac{2 \mathrm{Id}+\hat{U}_{\mathcal{P}}+\hat{U}_{\mathcal{P}}^{-1}}{4}\right) .
$$

It remains to compute $\left(\operatorname{det} \mathcal{L}_{R}\right)^{-2}$. To do so, we note that

$$
\operatorname{ker}\left(\mathcal{D}_{\left.\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}\right)} \equiv \operatorname{Im}(\operatorname{Id}-\mathcal{P}) \cap \operatorname{Im}\left(\Pi_{+}^{\sigma}\right)=\operatorname{Im}(\operatorname{Id}-\mathcal{P}) \cap \operatorname{Im}\left(\Pi_{+}\right)\right.
$$

since $\operatorname{Im}(\operatorname{Id}-\mathcal{P})$ and $\operatorname{Im}\left(\Pi_{+}^{\sigma}\right)$ have zero intersection in $\operatorname{ker}\left(D_{Y}\right)$ by definition of $\sigma$. Let $\left\{\psi_{\mu}\right\}$ be an orthonormal basis of $\operatorname{Im}(\operatorname{Id}-\mathcal{P}) \cap \operatorname{Im}\left(\Pi_{+}\right)$where $\psi_{\mu} \in E_{\mu}$. (Here, we recall that $D_{Y} \mathcal{P}=\mathcal{P} D_{Y}$, so $D_{Y}$ preserves $\operatorname{Im}(\operatorname{Id}-\mathcal{P})$ and obviously $\operatorname{Im}\left(\Pi_{+}\right)$as well. It follows that $D_{Y}$ can be diagonalized within the finite-dimensional space $\operatorname{Im}(\operatorname{Id}-\mathcal{P}) \cap \operatorname{Im}\left(\Pi_{+}\right)$. This elementary fact will be used quite often in the sequel.) Then $\phi_{\mu}=e^{-\mu u} \psi_{\mu} \in \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}\right)$ and

$$
\left\|\phi_{\mu}\right\|^{2}=\int_{-R}^{R} e^{-2 \mu u} d u=\frac{e^{2 \mu R}-e^{-2 \mu R}}{2 \mu}=\frac{\sinh (2 \mu R)}{\mu} .
$$

It follows that

$$
\mathcal{L}_{R}=\sum_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}(\mathrm{Id}-\mathcal{P}) \neq 0}} \frac{\phi_{\mu}(R)}{\left\|\phi_{\mu}\right\|} \otimes \frac{\phi_{\mu}(R)}{\left\|\phi_{\mu}\right\|}=\sum_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}(\mathrm{Id}-\mathcal{P}) \neq 0}} \frac{e^{-2 \mu R}}{\left\|\phi_{\mu}\right\|^{2}} \psi_{\mu} \otimes \psi_{\mu}
$$

Hence,

$$
\left(\operatorname{det} \mathcal{L}_{R}\right)^{-2}=\prod_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}(\operatorname{Id}-\mathcal{P}) \neq 0}}\left(\frac{e^{-2 \mu R}}{\left\|\phi_{\mu}\right\|^{2}}\right)^{-2}=\prod_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}(\mathrm{Id}-\mathcal{P}) \neq 0}} e^{4 \mu R} \frac{\sinh ^{2}(2 \mu R)}{\mu^{2}} .
$$

This completes the Gluing proof of Lemma 3.1.
Comparison proof of Lemma 3.1: We shall apply the comparison formula (2.5) to the pair $\left(\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{2}, \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{P}}^{2}\right)$. Here we regard Id $-\mathcal{P}$ as the Calderón projector at the boundary $\{R\} \times Y$ of the operator

$$
D: \operatorname{dom}(D) \rightarrow L^{2}\left(N_{R}, S\right),
$$

where

$$
\operatorname{dom}(D):=\left\{\phi \in H^{1}\left(N_{R}, S\right) \mid \mathcal{P}\left(\left.\phi\right|_{u=-R}\right)=0\right\} .
$$

Then $D_{\mathrm{Id}-\Pi_{+}^{\sigma}}=\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}$ and $D_{{c_{-}}}=\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{P}}$, so by the comparison formula (2.5),

$$
\frac{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{2}}{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{P}}^{2}}=(\operatorname{det} L)^{2} \cdot \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{U}_{\mathcal{P}}+\hat{U}_{\mathcal{P}}^{-1}}{4}\right),
$$

where we used that the $U$ in (2.5) for this situation is $\left(-\kappa_{\mathcal{P}}\right)\left(-\kappa_{\sigma}^{-1}\right)=\kappa_{\mathcal{P}} \kappa_{\sigma}^{-1}=: U_{\mathcal{P}}$, and where $L$ is the map over $\{R\} \times Y$ defined in (2.3):

$$
L=-P_{W} G \mathcal{R}^{-1} G P_{W}: W \rightarrow W
$$

with

$$
W:=\gamma_{R} \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{\sigma}\right) \equiv \operatorname{Im}(\operatorname{Id}-\mathcal{P}) \cap \operatorname{Im}\left(\Pi_{+}^{\sigma}\right)=\operatorname{Im}(\operatorname{Id}-\mathcal{P}) \cap \operatorname{Im}\left(\Pi_{+}\right) .
$$

To determine $(\operatorname{det} L)^{2}$, we need to find $\mathcal{R}$. An invertible extension $\tilde{D}$ of $D$ is just $G\left(\partial_{u}+D_{Y}\right)$ over $[-R, 2 R] \times Y$ with boundary condition $\mathcal{P}$ at $-R$ and $\operatorname{Id}-\mathcal{P}$ at $2 R$. Let $\left\{\psi_{\mu}\right\}$ be an orthonormal basis of $W=\operatorname{Im}(\operatorname{Id}-\mathcal{P}) \cap \operatorname{Im}\left(\Pi_{+}\right)$where $\psi_{\mu} \in E_{\mu}$, and for each such $\mu$, define $\varphi_{\mu}:=G \psi_{\mu} \in G W$. Since $G D_{Y}=-D_{Y} G$ and $\mathcal{P} G=G(\operatorname{Id}-\mathcal{P})$, it follows that $\varphi_{\mu}=G \psi_{\mu} \in E_{-\mu} \cap \operatorname{Im}(\mathcal{P})$. Using this, it is straightforward to check that

$$
\phi_{1}:=\frac{\sinh (\mu(u+R))}{\sinh (2 \mu R)} \varphi_{\mu} \quad \text { over } \quad[-R, R] \times Y,
$$

and

$$
\phi_{2}:=e^{\mu(u-R)} \varphi_{\mu} \quad \text { over } \quad[R, 2 R] \times Y,
$$

are continuous at $u=R$ with value $\varphi_{\mu}$ and satisfy $\widetilde{D}^{2} \phi_{i}=0, i=1,2$. Thus, by the definition of $\mathcal{R}$ in (2.4), we have

$$
\mathcal{R} \varphi_{\mu}:=\left.\partial_{u} \phi_{1}\right|_{u=R}-\left.\partial_{u} \phi_{2}\right|_{u=R}=\left(\mu \frac{\cosh (2 \mu R)}{\sinh (2 \mu R)}-\mu\right) \varphi_{\mu}=\frac{\mu e^{-2 \mu R}}{\sinh (2 \mu R)} \varphi_{\mu} .
$$

Therefore,

$$
-G \mathcal{R} G \psi_{\mu}=-G \mathcal{R} \varphi_{\mu}=-G \frac{\mu e^{-2 \mu R}}{\sinh (2 \mu R)} \varphi_{\mu}=\frac{\mu e^{-2 \mu R}}{\sinh (2 \mu R)} \psi_{\mu}
$$

It follows that

$$
L=-P_{W} G \mathcal{R}^{-1} G P_{W}=\sum_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}(\operatorname{Id}-\mathcal{P}) \neq 0}} \frac{e^{2 \mu R} \sinh (2 \mu R)}{\mu} \psi_{\mu} \otimes \psi_{\mu}
$$

Hence,

$$
(\operatorname{det} L)^{2}=\prod_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}(\mathrm{Id}-\mathcal{P}) \neq 0}} e^{4 \mu R} \frac{\sinh ^{2}(2 \mu R)}{\mu^{2}},
$$

which completes the Comparison proof of Lemma 3.1.
Next, we compute a related $\zeta$-determinant ratio.
Lemma 3.2: With the notations above, the following equality holds:

$$
\left.\frac{\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma}, \mathrm{Id}-\mathcal{P}}^{2}}{\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{2}} \prod_{\substack{\mu>0 \\ E_{\mu} \cap \mathrm{Im}(\mathrm{Id}-\mathcal{P}) \neq 0}} \frac{e^{-4 \mu R} \sinh ^{2}(2 \mu R)}{\mu^{2}}\right) \cdot \operatorname{det}_{F}\left(\frac{2 \mathrm{Id}+\hat{U}_{\mathcal{P}}+\hat{U}_{\mathcal{P}}^{-1}}{4}\right)
$$

Proof: Observe that $\operatorname{Id}-\Pi_{+}^{\sigma}$ is the Calderón projector at $\{R\} \times Y$ of the operator

$$
D: \operatorname{dom}(D) \rightarrow L^{2}\left(N_{R}, S\right)
$$

where

$$
\operatorname{dom}(D):=\left\{\phi \in H^{1}\left(N_{R}, S\right) \mid \Pi_{+}^{\sigma}\left(\left.\phi\right|_{u=-R}\right)=0\right\} .
$$

Then $D_{\text {Id }-\mathcal{P}}=\mathcal{D}_{\Pi_{+}^{\sigma} \text { Id }-\mathcal{P}}$ and $D_{\mathcal{C}_{-}}=\mathcal{D}_{\Pi_{+}^{\sigma}, \text { Id }-\Pi_{+}^{\sigma}}$, so the comparison formula (2.5) applied to this situation gives us

$$
\frac{\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma}, \mathrm{Id}-\mathcal{P}}^{2}}{\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{2}}=(\operatorname{det} L)^{2} \cdot \operatorname{det}_{F}\left(\frac{2 \mathrm{Id}+\hat{U}_{\mathcal{P}}^{-1}+\hat{U}_{\mathcal{P}}}{4}\right),
$$

where we used that the $U$ in (2.5) for this situation is $\left(-\kappa_{\sigma}\right)\left(-\kappa_{\mathcal{P}}^{-1}\right)=\kappa_{\sigma} \kappa_{\mathcal{P}}^{-1}=\left(\kappa_{\mathcal{P}} \kappa_{\sigma}^{-1}\right)^{-1}=U_{\mathcal{P}}^{-1}$, and where $L$ is the map over $\{R\} \times Y$ defined in (2.3):

$$
L:=-P_{W} G \mathcal{R}^{-1} G P_{W}: W \rightarrow W
$$

with

$$
W:=\gamma_{R} \operatorname{ker}\left(\mathcal{D}_{\Pi_{+}^{\sigma}} \mathrm{Id}-\mathcal{P}\right) \equiv \operatorname{Im}\left(\operatorname{Id}-\Pi_{+}^{\sigma}\right) \cap \operatorname{Im}(\mathcal{P})=\operatorname{Im}(\mathcal{P}) \cap \operatorname{Im}\left(\Pi_{-}\right)
$$

To compute $(\operatorname{det} L)^{2}$ for this example, we proceed in much the same way as for the Comparison proof of Lemma 3.1. An invertible extension $\tilde{D}$ of $D$ is just $G\left(\partial_{u}+D_{Y}\right)$ over $[-R, 2 R] \times Y$ with boundary condition $\Pi_{+}^{\sigma}$ at $-R$ and $\operatorname{Id}-\Pi_{+}^{\sigma}$ at $2 R$. Let $\left\{\psi_{\nu}\right\}$ be an orthonormal basis of $W$ $=\operatorname{Im}(\mathcal{P}) \cap \operatorname{Im}\left(\Pi_{-}\right)$where $\psi_{\nu} \in E_{\nu}$. It is important to note that, in contrast to the Comparison proof of Lemma 3.1, the $\nu$ 's are negative (rather than positive) and that $\psi_{\nu} \in E_{\nu} \cap \operatorname{Im}(\mathcal{P})$ [rather than $\operatorname{Im}(\operatorname{Id}-\mathcal{P})]$. Following the Comparison proof of Lemma 3.1 almost verbatim, it is straightforward to check that

$$
L=-P_{W} G \mathcal{R}^{-1} G P_{W}=\sum_{\substack{\nu<0 \\ E_{\nu} \cap \operatorname{Im}(\mathcal{P}) \neq 0}} \frac{e^{2 \nu R} \sinh (2 \nu R)}{\nu} \psi_{\nu} \otimes \psi_{\nu}
$$

so

$$
(\operatorname{det} L)^{2}=\prod_{\substack{\nu<0 \\ E_{\nu} \cap \operatorname{Im}(\mathcal{P}) \neq 0}} e^{4 \nu R} \frac{\sinh ^{2}(2 \nu R)}{\nu^{2}}
$$

Finally, using that $G$ maps $E_{\nu} \cap \operatorname{Im}(\mathcal{P})$ isomorphically onto $E_{-\nu} \cap \operatorname{Im}(\operatorname{Id}-\mathcal{P})$, where we used that $G D_{Y}=-D_{Y} G$ and $\mathcal{P} G=G(\mathrm{Id}-\mathcal{P})$, we finally get

$$
(\operatorname{det} L)^{2}=\prod_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}(\operatorname{Id}-\mathcal{P}) \neq 0}} e^{4(-\mu) R} \frac{\sinh ^{2}(-2 \mu R)}{(-\mu)^{2}}=\prod_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}(\mathrm{Id}-\mathcal{P}) \neq 0}} e^{-4 \mu R} \frac{\sinh ^{2}(2 \mu R)}{\mu^{2}},
$$

which completes our proof.
We are now ready to prove the following
Lemma 3.3: With the notations above, the following equality holds:

$$
\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{P}}^{2}=e^{2 C R_{2}^{\zeta} b_{Y}^{2}(0)+h_{Y}} \cdot\left(\prod_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}(\mathrm{Id}-\mathcal{P}) \neq 0}} e^{-8 \mu R}\right),
$$

where $C=\left(\Gamma(s)^{-1} \zeta_{D_{Y}^{2}}(s-1 / 2)\right)^{\prime}(0)$ with $\zeta_{D_{Y}^{2}}(s)$ the $\zeta$-function of $D_{Y}^{2}$ and $h_{Y}=\operatorname{dim} \operatorname{ker}\left(D_{Y}\right)$.

Proof: Solving for $\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{P}}^{2}$ in Lemma 3.1, we obtain

$$
\begin{equation*}
\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{P}}^{2}=\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{2} \cdot\left(\prod_{\substack{\mu>0 \\ E_{\mu} \cap \mathrm{Im}(\mathrm{Id}-\mathcal{P}) \neq 0}} \frac{e^{-4 \mu R} \mu^{2}}{\sinh ^{2}(2 \mu R)}\right) \cdot \operatorname{det}_{F}\left(\frac{2 \mathrm{Id}+\hat{U}_{\mathcal{P}}+\hat{U}_{\mathcal{P}}^{-1}}{4}\right)^{-1} \tag{3.3}
\end{equation*}
$$

On the other hand, from Lemma 3.2, we know that

$$
\left.\frac{\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}}^{2} \mathrm{Id}-\mathcal{P}}{\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}}^{2}} \prod_{\substack{\mu \mathrm{Id}-\Pi_{+}^{\sigma}}} \frac{e^{-4 \mu R} \sinh ^{2}(2 \mu R)}{\mu^{2}}\right) \cdot \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{U}_{\mathcal{P}}+\hat{U}_{\mathcal{P}}^{-1}}{4}\right)
$$

and from the main result of Ref. 11 [see formula (1.7)], we also have

$$
\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{2}=e^{2 C R_{2}{ }^{\zeta} D_{Y}^{2}(0)+h_{Y}}
$$

Hence,

$$
\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma}, \mathrm{Id}-\mathcal{P}}^{2}=e^{2 C R} 2^{\zeta D_{Y}^{2}(0)+h_{Y}}\left(\prod_{\substack{\mu>0 \\ E_{\mu} \cap \mathrm{Im}(\mathrm{Id}-\mathcal{P}) \neq 0}} \frac{e^{-4 \mu R} \sinh ^{2}(2 \mu R)}{\mu^{2}}\right) \cdot \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{U}_{\mathcal{P}}+\hat{U}_{\mathcal{P}}^{-1}}{4}\right)
$$

Substituting this expression into (3.3), using that

$$
\begin{equation*}
\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\Pi_{+}^{\sigma}}^{2}=\operatorname{det}_{\zeta} \mathcal{D}_{\Pi_{+}^{\sigma}, \mathrm{Id}-\mathcal{P}}^{2}, \tag{3.4}
\end{equation*}
$$

concludes the proof of our result once we show that (3.4) holds. In fact, (3.4) holds in the more general setting: $\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{Q}}^{2}=\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{Q}, \mathrm{Id}-\mathcal{P}}^{2}$ for all $(\mathcal{P}, \operatorname{Id}-\mathcal{Q}) \in \operatorname{Gr}_{\infty}^{*}\left(D_{Y}\right)$. To prove this, we simply observe that

$$
G\left(\partial_{u}+D_{Y}\right) \phi=\lambda \phi, \quad \mathcal{P} \phi(-R)=0, \quad(\operatorname{Id}-\mathcal{Q}) \phi(R)=0
$$

if and only if $\psi(u):=G \phi(-u)$ satisfies

$$
G\left(\partial_{u}+D_{Y}\right) \psi=-\lambda \psi, \quad(\mathrm{Id}-\mathcal{P}) \psi(R)=0, \quad \mathcal{Q} \psi(-R)=0,
$$

where we used that $G^{2}=-\mathrm{Id}, G D_{Y}=-D_{Y} G, \mathcal{P} G=G(\mathrm{Id}-\mathcal{P})$, and $\mathcal{Q} G=G(\mathrm{Id}-\mathcal{Q})$. It follows that $\operatorname{spec}\left(\mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{Q}}^{2}\right)=\operatorname{spec}\left(\mathcal{D}_{\mathcal{Q}, \mathrm{Id}-\mathcal{P}}^{2}\right)$, which implies that $\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}, \mathrm{Id}-\mathcal{Q}}^{2}=\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{Q}, \mathrm{Id}-\mathcal{P}}^{2}$.

## IV. PROOF OF THEOREM 1.1

As we already mentioned, in order to demonstrate the effectiveness of the gluing formula (2.2) and the comparison formula (2.5), we shall give separate proofs of Theorem 1.1 exploiting both formulas.

Gluing proof of Theorem 1.1: Recall that the operator $\mathcal{D}_{P}$ as defined in (1.4) written using the notation (3.1) is

$$
\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}:=\mathcal{D}: \operatorname{dom}\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}\right) \rightarrow L^{2}\left(N_{R}, S\right),
$$

where

$$
\operatorname{dom}\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}\right):=\left\{\phi \in H^{1}\left(N_{R}, S\right) \mid \mathcal{P}_{1}\left(\left.\phi\right|_{u=-R}\right)=0, \quad \mathcal{P}_{2}\left(\left.\phi\right|_{u=R}\right)=0\right\} .
$$

Let us decompose $N_{R}$ into two parts $[-R, 0] \times Y$ and $[0, R] \times Y$ as shown in Fig. 2. Then the restrictions of $\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}$ over the decomposed parts define two Dirac type operators with boundary


FIG. 2. Cutting $N_{R}$ at 0 into two pieces.
conditions given by $\mathcal{P}_{1}$ at $\{-R\} \times Y$ and $\mathcal{P}_{2}$ at $\{R\} \times Y$. It is easy to check that the Calderón projections of these operators are, respectively, just Id $-\mathcal{P}_{1}$ and Id $-\mathcal{P}_{2}$ over $\{0\} \times Y$. We denote the operators over the decomposed parts $[-R, 0] \times Y$ and $[0, R] \times Y$ with the new boundary conditions given by these Calderón projectors over $\{0\} \times Y$ by $D_{\mathcal{P}_{1}, \text { Id }-\mathcal{P}_{1}}$ and $D_{\text {Id }-\mathcal{P}_{2}, \mathcal{P}_{2}}$, respectively. Now applying the gluing formula (2.2) to this situation, we obtain

$$
\begin{equation*}
\frac{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}^{2}}{\operatorname{det}_{\zeta} D_{\mathcal{P}_{1}, \mathrm{Id}-\mathcal{P}_{1}}^{2} \cdot \operatorname{det}_{\zeta} D_{\mathrm{Id}-\mathcal{P}_{2}, \mathcal{P}_{2}}^{2}}=2^{-\zeta_{D_{Y}^{2}}^{2}(0)-h_{Y}}(\operatorname{det} \mathcal{L})^{-2} \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{U}_{P}+\hat{U}_{P}^{-1}}{4}\right), \tag{4.1}
\end{equation*}
$$

where the $\hat{\mathcal{U}}$ in (2.2) for this situation is the restriction of $-\left(-\kappa_{1}\right)\left(-\kappa_{2}^{-1}\right)=-\kappa_{1} \kappa_{2}^{-1}=$ : $U_{P}$ over $L^{2}\left(Y, S^{-}\right)$to the orthogonal complement of its $(-1)$-eigenspace, noting that $-\kappa_{i}$ is the isometry corresponding to $\operatorname{Id}-\mathcal{P}_{i}$ for $i=1,2$, and where

$$
\begin{equation*}
\mathcal{L}:=\sum_{k=1}^{h_{P}} \gamma_{0} U_{k} \otimes \gamma_{0} U_{k}: \gamma_{0} \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}\right) \rightarrow \gamma_{0} \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}\right), \tag{4.2}
\end{equation*}
$$

where $h_{P}=\operatorname{dim} \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}\right), \gamma_{0}$ is the restriction map from $N_{R}$ to $\{0\} \times Y$, and $\left\{U_{k}\right\}$ is an orthonormal basis of $\operatorname{ker}\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}\right)$. By Lemma 3.3 (here we need to replace $R$ with $R / 2$ since the lengths of $[-R, 0]$ and $[0, R]$ are half that of $[-R, R]$, which is the interval considered in Lemma 3.3),

$$
\operatorname{det}_{\zeta} D_{\mathcal{P}_{1}, \mathrm{Id}-\mathcal{P}_{1}}^{2}=e^{C R_{2} \zeta_{D_{Y}^{2}}^{2}(0)+h_{Y}} \cdot\left(\prod_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}\left(\mathrm{Id}-\mathcal{P}_{1}\right) \neq 0}} e^{-4 \mu R}\right),
$$

and

Now to complete the Gluing proof of Theorem 1.1, it remains to compute $(\operatorname{det} \mathcal{L})^{-2}$ over $\gamma_{0} \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}\right) \equiv \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{2}\right)$. To do so, we note that

$$
\operatorname{ker}\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}\right) \equiv \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{2}\right):=W
$$

Let $\left\{\psi_{\mu}\right\}$ be an orthonormal basis of $W$ where $\psi_{\mu} \in E_{\mu}$. Then $\phi_{\mu}:=e^{-\mu u} \psi_{\mu} \in \operatorname{ker}\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}\right)$ and

$$
\left\|\phi_{\mu}\right\|^{2}=\int_{-R}^{R} e^{-2 \mu u} d u=\frac{\sinh (2 \mu R)}{\mu}
$$

where if $\mu=0$, then we replace $[\sinh (2 \mu R)] / \mu$ by its limit as $\mu \rightarrow 0$, that is, $2 R$. It follows that

$$
\mathcal{L}=\sum_{\substack{\mu \in \operatorname{spec}\left(D_{Y}\right) \\ E_{\mu} \cap W \neq 0}} \frac{\phi_{\mu}(0)}{\left\|\phi_{\mu}\right\|} \otimes \frac{\phi_{\mu}(0)}{\left\|\phi_{\mu}\right\|}=\sum_{\substack{\mu \in \operatorname{spec}\left(D_{Y}\right) \\ E_{\mu} \cap W \neq 0}} \frac{1}{\left\|\phi_{\mu}\right\|^{2}} \psi_{\mu} \otimes \psi_{\mu}
$$

Hence,

$$
(\operatorname{det} \mathcal{L})^{-2}=\prod_{\substack{\mu \in \operatorname{spec}\left(D_{y}\right) \\ E_{\mu} \cap W \neq 0}}\left(\frac{1}{\left\|\phi_{\mu}\right\|^{2}}\right)^{-2}=\prod_{\substack{\mu \in \operatorname{spec}\left(D_{y}\right) \\ E_{\mu} \cap W \neq 0}} \frac{\sinh ^{2}(2 \mu R)}{\mu^{2}} .
$$

This completes the Gluing proof of Theorem 1.1.
Comparison proof of Theorem 1.1: We now prove Theorem 1.1 using the comparison formula (2.5) applied to the pair $\left(\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}^{2}, \mathcal{D}_{\mathcal{P}_{1}, \mathrm{Id}-\mathcal{P}_{1}}^{2}\right)$. Here we regard Id $-\mathcal{P}_{1}$ as the Calderón projector at the boundary $\{R\} \times Y$ of the operator

$$
D: \operatorname{dom}(D) \rightarrow L^{2}\left(N_{R}, S\right)
$$

where

$$
\operatorname{dom}(D):=\left\{\phi \in H^{1}\left(N_{R}, S\right) \mid \mathcal{P}\left(\left.\phi\right|_{u=-R}\right)=0\right\} .
$$

Then $D_{\mathcal{P}_{2}}=\mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}$ and $D_{c_{-}}=\mathcal{D}_{\mathcal{P}_{1}, \mathrm{Id}-\mathcal{P}_{1}}$, so applying the comparison formula (2.5) to this situation, we obtain

$$
\begin{equation*}
\frac{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}^{2}}{\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}_{1}, \mathrm{Id}-\mathcal{P}_{1}}^{2}}=(\operatorname{det} L)^{2} \cdot \operatorname{det}_{F}\left(\frac{2 \mathrm{Id}+\hat{U}_{P}+\hat{U}_{P}^{-1}}{4}\right) \tag{4.3}
\end{equation*}
$$

where we used that the $U$ in (2.5) for this situation is $\left(-\kappa_{1}\right) \kappa_{2}^{-1}=$ : $U_{P}$ noting that $-\kappa_{1}$ corresponds to Id $-\mathcal{P}_{1}$, and where $L$ is the operator defined in (2.3) for this situation, which we will investigate in detail soon. Now, by Lemma 3.3,

$$
\left.\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}_{1}, \mathrm{Id}-\mathcal{P}_{1}}^{2}=e^{2 C R_{2}^{\zeta D_{Y}^{2}(0)+h_{Y}}} \cdot \prod_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \neq 0}} e^{-8 \mu R}\right),
$$

therefore by (4.3),

$$
\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}^{2}=e^{2 C R_{2}{ }^{\zeta D_{Y}^{2}(0)+h_{Y}}}\left(\prod_{\substack{\mu>0 \\ E_{\mu} \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \neq 0}} e^{-8 \mu R}\right) \cdot(\operatorname{det} L)^{2} \cdot \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{U}_{P}+\hat{U}_{P}^{-1}}{4}\right) .
$$

To compute det $L$, we use almost the exact same argument found in Lemma 3.2 to show that with $W:=\operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{2}\right)$,

$$
(\operatorname{det} L)^{2}=\prod_{\substack{\mu \in \operatorname{spec}\left(D_{Y}\right) \\ E_{\mu} \cap W \neq 0}} e^{4 \mu R} \frac{\sinh ^{2}(2 \mu R)}{\mu^{2}},
$$

where in the product, when $\mu=0$ we replace $\left[\sinh ^{2}(2 \mu R)\right] / \mu^{2}$ by its limit as $\mu \rightarrow 0$, that is, $(2 R)^{2}$. Therefore,

$$
\begin{align*}
\operatorname{det}_{\zeta} \mathcal{D}_{\mathcal{P}_{1}, \mathcal{P}_{2}}^{2}= & e^{2 C R_{2}{ }_{2} D_{Y}^{2}(0)+h_{Y}}\left(\prod_{\substack{\mu>0 \\
E_{\mu} \cap \operatorname{Im}\left(\mathrm{Id}-\mathcal{P}_{1}\right) \neq 0}} e^{-8 \mu R}\right) \cdot\left(\prod_{\substack{\mu \in \operatorname{spec}\left(D_{Y}\right) \\
E_{\mu} \cap W \neq 0}} e^{4 \mu R}\right) \\
& \cdot\left(\prod_{\substack{\mu \in \operatorname{spec}\left(D_{Y}\right) \\
E_{\mu} \cap W \neq 0}} \frac{\sinh ^{2}(2 \mu R)}{\mu^{2}}\right) \cdot \operatorname{det}_{F}\left(\frac{2 \operatorname{Id}+\hat{U}_{P}+\hat{U}_{P}^{-1}}{4}\right) . \tag{4.4}
\end{align*}
$$

Since $W=\operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{2}\right)$, we have


Since $D_{Y} G=-G D_{Y}$ and $\mathcal{P}_{i} G=G\left(\operatorname{Id}-\mathcal{P}_{i}\right)$, we have

$$
E_{\mu} \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{2}\right) \neq 0 \Leftrightarrow E_{-\mu} \cap \operatorname{Im}\left(\mathcal{P}_{1}\right) \cap \operatorname{Im} \mathcal{P}_{2} \neq 0
$$

as $G$ maps the first space isomorphically onto the second space, therefore

$$
\left.\begin{array}{rl}
\prod_{\substack{\mu \in \operatorname{spec}\left(D_{Y}\right) \\
E_{\mu} \cap W \neq 0}} e^{4 \mu R}= & \prod_{\substack{\mu>0 \\
E_{\mu} \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \neq 0}} e^{4 \mu R} \\
& \cdot\left(\begin{array}{cc}
\prod_{\mu>0}^{\mu>0} \\
\sum_{\mu} \cap \operatorname{Im}\left(\operatorname{Id}-\mathcal{P}_{1}\right) \cap \operatorname{Im}\left(\mathcal{P}_{2}\right) \neq 0
\end{array}\right. \\
\prod_{E_{\mu} \cap \operatorname{Im}\left(\mathcal{P}_{1}\right) \cap \operatorname{Im}\left(\mathcal{P}_{2}\right) \neq 0}
\end{array}\right) \cdot\left(e^{-4 \mu R}\right)
$$

Putting this expression into (4.4) completes the Comparison proof of Theorem 1.1.
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