## PROOFS ABOUT THE ORLIK–SOLOMON ALGEBRA A OF A MATROID

Throughout, M is a simple matroid. The ground set  $\mathcal E$  is linearly ordered. A broken circuit is a set  $C \setminus \max C$ , where C is a circuit and  $\max C$  denotes the largest element of C in the ground-set linear ordering. Recall that we defined  $NBC(\mathcal{M})$  to be the class of no-broken-circuit sets (NBC-sets, not sometimes called "peacock sets"), i.e., sets  $S \subseteq \mathcal{E}$  that do not contain any broken circuit. We defined  $NBC_p = \{S \in \text{NBC} : |S| = p\}.$ 

In the algebras, when I write for instance  $a<sub>S</sub>$  where S is a sequence of matroid elements, I will assume they are in any order, i.e., I don't distinguish between different orderings of S. The specific ordering can't affect anything but the sign, so all orderings give the same basis element for practical purposes. I can make this formal: assume the chosen  $S$  is the one in the ground-set linear order.

## **Theorem 0.1.** The elements  $a_N$  for  $N \in \text{NBC}_p(\mathcal{M})$  form a basis for  $A_p$ .

*Proof.* Consider a circuit  $C = \{h, h_1, h_2, \ldots, h_k \text{ where } h = \max C \text{ (by choice of notation).} \}$ Then  $\partial e_C = e_{C\setminus h} + \sum_i \pm e_{C\setminus h_i}$ , and  $\partial e_C \in I$ , so  $a_{C\setminus h} = \sum_i \pm a_{C\setminus h_i}$ .  $C \setminus h$  is a broken circuit; this shows that to generate  $A$  as a vector space we don't need any broken circuit terms. Since I is an ideal, we could have said we don't need any terms  $a_S a_{C\hbar}$ , since  $a_S a_{C\setminus h} = sum_i \pm a_S a_{C\setminus h_i}$ . That is, we only need no-broken-circuit terms; in particular, A is spanned by  $a_N$  for  $N \in \text{NBC}_p(\mathcal{M})$ .

To show this set is a basis we count. We proved dim  $A_p = |w_p(\mathcal{M})|$ . Brylawski proved  $|\text{NBC}_p(\mathcal{M})| = |w_p(\mathcal{M})|$ . QED

Convenient notation:  $\mathcal{M}' = \mathcal{M} \setminus h$ ,  $\mathcal{M}'' = \mathcal{M}/h$ .  $A' = A(\mathcal{M}')$ , etc.

<span id="page-0-0"></span>**Theorem 0.2.** Assume h is not a coloop in M. The sequence  $0 \stackrel{i}{\rightarrow} A'_{p} \rightarrow A_{p} \stackrel{j}{\rightarrow} A''_{p-1} \rightarrow 0$ is exact.

*Proof.* Obviously,  $ji = 0$ , i.e., Ker  $j \subset \text{Im } i$ . We know the dimensions:

 $\dim A_p = |w_p(\mathcal{M})|, \quad \dim A'_p = |w_p(\mathcal{M}')|, \quad \dim A''_{p-1} = |w_{p-1}(\mathcal{M}'')|.$ 

Standard matroid theory (by Rota): if h is neither a loop nor a coloop in M, then  $p_{\mathcal{M}}(\lambda) =$  $p_{\mathcal{M}'}(\lambda) - p_{\mathcal{M}''}(\lambda)$ . From this, remembering that  $w_p(\mathcal{M})$  is the coefficient of  $\lambda^{r(\mathcal{M})-p}$ , we can deduce that  $w_p(\mathcal{M}) = w_p(\mathcal{M}') - w_{p-1}(\mathcal{M}'')$ . Rota's sign theorem then implies that  $|w_p(\mathcal{M})| =$  $|w_p(\mathcal{M}')| + |w_{p-1}(\mathcal{M}'')|$ . Restating this in terms of A's,

$$
\dim A_p = \dim A'_p + \dim A''_{p-1},
$$

which gives dim Ker  $j = \dim A_p - \dim A_{p-1}'' = \dim A_p' = \dim \text{Im } i$  and voila!

These proofs show how important is Orlik's Proposition 3.20, the one that says dim  $A_X =$  $(-1)^{r(X)}\mu(\hat{0},X)$ , from which we inferred that dim  $A_p = |w_p|$ .

Corollary 0.3. The short exact sequence of Theorem [0.2](#page-0-0) splits.

*Proof.* From algebra, since the outer terms are free  $\mathbb{Z}$ -modules.  $\Box$ 

(Thanks to Chris for timely reminders about the dimension law and splitting.)

— Submitted for your approval by Tom; 26, 29 Feb 2020